

Mid-year Report: Lagrangian Analysis of Two- and Three-Dimensional Oceanic Flows from Eulerian Velocity Data

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Project Overview

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- Project goal: Create tools for Lagrangian analysis of oceanic flow (2D or 3D), given only velocity data on spatio-temporal grid (e.g. from model output)
- Track vast network of particles through flow, use trajectories to bring out underlying Lagrangian coherent structures, as well as stable and unstable manifolds separating these structures
- Test tools on velocity output from Chesapeake Bay model

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Example: Kuroshio current, northwest Pacific Ocean

- Red indicates faster-moving regions, blue slower
- Thin yellow lines represent stable and unstable manifolds

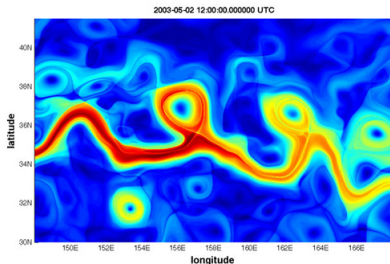


Figure: Coherent structures in the Kuroshio current [1]

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- Velocity field $\mathbf{u}(x_i, y_j, z_k, t_l) = (u, v, w)$ given at discrete points in space and time
- Want to compute particle trajectory $\mathbf{X}(\mathbf{X}_0, t)$ given initial position $\mathbf{X}_0 \in \mathbb{R}^2$ or \mathbb{R}^3
- Two numerical tasks: *interpolation* and *time integration*
 - *Interpolation*: Particle velocities $\mathbf{u}(\mathbf{X}(\mathbf{X}_0, t))$ must be interpolated from grid velocities $\mathbf{u}_{i,j,k,l}$
 - *Integration*: Velocity must be integrated in time to obtain position as time evolves:

$$\mathbf{X}(\mathbf{X}_0, t) = \int_0^t \mathbf{u}(\mathbf{X}(\mathbf{X}_0, t'), t') dt'$$

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- Given trajectories, two main tools for Lagrangian analysis: *M-function* and *largest finite-time Lyapunov exponent (FTLE)*
- *M-function*: distance traveled by particle within some fixed time interval (e.g. τ forward and backward from current time):

$$M_{\mathbf{u},\tau}(\mathbf{X}_0, t) = \int_{t-\tau}^{t+\tau} \left(\sum_{i=1}^{2 \text{ or } 3} \left(\frac{dX_i(\mathbf{X}_0, t')}{dt'} \right)^2 \right)^{\frac{1}{2}} dt'$$

- Coloring by *M-function* brings out boundaries between coherent structures moving at different speeds

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- Given trajectories, two main tools for Lagrangian analysis:
M-function and *largest finite-time Lyapunov exponent (FTLE)*
- *FTLE*: exponential growth rate along maximal growth axis of an infinitesimal parcel of fluid:

$$\lambda(t) = \frac{1}{2t} \ln \left(\rho(L^T L) \right)$$

where ρ denotes the spectral radius and $L(t) = \frac{\partial \mathbf{X}(\mathbf{X}_0, t)}{\partial \mathbf{X}_0}$.

- Coloring by FTLE highlights bifurcation regions in flow

Implementation Basics

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- Store particle positions in column vectors X_p , Y_p , Z_p
- Basic principle: vectorize all operations (all particles at once rather than looping through them)
- Input velocity data in Arakawa C-grid format (u given at east and west sides of grid box, v at north and south, w at top and bottom):

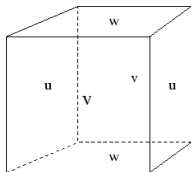


Figure: Arakawa c-grid box [5]

2D Interpolation Algorithms

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- Compare two methods for 2D spatial interpolation: piecewise *bilinear* and piecewise *bicubic* functions (splines)
 - Bilinear: Within each grid box, fit function of form $p(x, y) = \sum_{i,j=0}^1 a_{ij}x^i y^j = a_{00} + a_{10}x + a_{01}y + a_{11}xy$ to four corner values of u
 - Bicubic: Fit function of form $p(x, y) = \sum_{i,j=0}^3 a_{ij}x^i y^j$ to four corner values of u and estimates of its derivatives u_x , u_y , and u_{xy}
- Bicubic should be slower but more accurate

Bilinear Interpolation

- Approximate $u(x, y)$ on $[x_i, x_{i+1}] \times [y_j, y_{j+1}]$ by

$$p(x, y) = a_{00} + a_{10}x + a_{01}y + a_{11}xy$$

- Explicit formula:

$$p(x, y) = w_y (w_x u_{i,j} + (1 - w_x) u_{i+1,j}) \\ + (1 - w_y) (w_x u_{i,j+1} + (1 - w_x) u_{i+1,j+1})$$

with weights

$$w_x = \frac{x_{i+1} - x}{x_{i+1} - x_i}$$

$$w_y = \frac{y_{j+1} - y}{y_{j+1} - y_j}$$

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- Approximate $u(x, y)$ on $[0, 1]^2$ by

$$p(x, y) = \sum_{i,j=0}^3 a_{ij} x^i y^j$$

- Find a_{ij} 's by matching function to known values $u(0, 0)$, $u(0, 1)$, $u(1, 0)$, $u(1, 1)$ and finite difference approximations for partial derivatives u_x , u_y , and u_{xy} at corners
- Sixteen parameters, sixteen unknowns
- Scale inputs from grid box to $[0, 1]^2$, solve, scale back

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- Second-order centered finite difference approximations for *interior* grid points:

$$(u_x)_{i,j} \approx \frac{1}{2\Delta x} \begin{bmatrix} -1 & 0 & 1 \end{bmatrix} \cdot \begin{bmatrix} u_{i-1,j} \\ u_{i,j} \\ u_{i+1,j} \end{bmatrix} = \frac{u_{i+1,j} - u_{i-1,j}}{2\Delta x}$$

$$(u_y)_{i,j} \approx \begin{bmatrix} u_{i,j-1} & u_{i,j} & u_{i,j+1} \end{bmatrix} \cdot \frac{1}{2\Delta y} \begin{bmatrix} -1 \\ 0 \\ 1 \end{bmatrix} = \frac{u_{i,j+1} - u_{i,j-1}}{2\Delta y}$$

$$(u_{xy})_{i,j} \approx \frac{1}{2\Delta x} \begin{bmatrix} -1 & 0 & 1 \end{bmatrix} \cdot \begin{bmatrix} u_{i-1,j-1} & u_{i-1,j} & u_{i-1,j+1} \\ u_{i,j-1} & u_{i,j} & u_{i,j+1} \\ u_{i+1,j-1} & u_{i+1,j} & u_{i+1,j+1} \end{bmatrix} \cdot \frac{1}{2\Delta y} \begin{bmatrix} -1 \\ 0 \\ 1 \end{bmatrix} \\ = \frac{u_{i+1,j+1} - u_{i+1,j-1} - u_{i-1,j+1} + u_{i-1,j-1}}{4\Delta x \Delta y}$$

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- Second-order one-sided finite difference approximations for *boundary* grid points:

$$(u_x)_{1,j} \approx \frac{1}{2\Delta x} [-3 \quad 4 \quad -1] \cdot \begin{bmatrix} u_{1,j} \\ u_{2,j} \\ u_{3,j} \end{bmatrix}$$

$$(u_{xy})_{1,j} \approx \frac{1}{2\Delta x} [-3 \quad 4 \quad -1] \cdot \begin{bmatrix} u_{1,j-1} & u_{1,j} & u_{1,j+1} \\ u_{2,j-1} & u_{2,j} & u_{2,j+1} \\ u_{3,j-1} & u_{3,j} & u_{3,j+1} \end{bmatrix} \cdot \frac{1}{2\Delta y} \begin{bmatrix} -1 \\ 0 \\ 1 \end{bmatrix}$$

$$(u_{xy})_{1,1} \approx \frac{1}{2\Delta x} [-3 \quad 4 \quad -1] \cdot \begin{bmatrix} u_{1,1} & u_{1,2} & u_{1,3} \\ u_{2,1} & u_{2,2} & u_{2,3} \\ u_{3,1} & u_{3,2} & u_{3,3} \end{bmatrix} \cdot \frac{1}{2\Delta y} \begin{bmatrix} -3 \\ 4 \\ -1 \end{bmatrix}$$
$$= \frac{9u_{1,1} - 12u_{1,2} + 3u_{1,3} - 12u_{2,1} + 16u_{2,2} - 4u_{2,3} + 3u_{3,1} - 4u_{3,2} + u_{3,3}}{4\Delta x \Delta y}$$

...and likewise for other boundary points

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- Determine a_{ij} from these values via

$$\begin{bmatrix} a_{00} \\ a_{10} \\ a_{20} \\ a_{30} \\ a_{01} \\ a_{11} \\ a_{21} \\ a_{31} \\ a_{02} \\ a_{12} \\ a_{22} \\ a_{32} \\ a_{03} \\ a_{13} \\ a_{23} \\ a_{33} \end{bmatrix} = \begin{bmatrix} 1 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 1 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\ -3 & 3 & 0 & 0 & -2 & -1 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\ 2 & -2 & 0 & 0 & 1 & 1 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 1 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 1 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & -3 & 3 & 0 & 0 & -2 & -1 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 2 & -2 & 0 & 0 & 1 & 1 & 0 & 0 \\ -3 & 0 & 0 & 3 & 0 & 0 & 0 & 0 & -2 & 0 & -1 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & -3 & 0 & 3 & 0 & 0 & 0 & 0 & -2 & 0 & -1 & 0 \\ 9 & -9 & -9 & 9 & 6 & 3 & -6 & -3 & 6 & -6 & 3 & -3 & 4 & 2 & 2 & 1 \\ -6 & 6 & 6 & -6 & -3 & -3 & 3 & 3 & -4 & 4 & -2 & 2 & -2 & -2 & -1 & -1 \\ 2 & 0 & -2 & 0 & 0 & 0 & 0 & 0 & 1 & 0 & 1 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 2 & 0 & -2 & 0 & 0 & 0 & 0 & 1 & 0 & 1 & 0 & 0 \\ -6 & 6 & 6 & -6 & -4 & -2 & 4 & 2 & -3 & 3 & -3 & 3 & -2 & -1 & -2 & -1 \\ 4 & -4 & -4 & 4 & 2 & 2 & -2 & -2 & 2 & -2 & 2 & -2 & 1 & 1 & 1 & 1 \end{bmatrix} \begin{bmatrix} u(0, 0) \\ u(1, 0) \\ u(0, 1) \\ u(1, 1) \\ u_x(0, 0) \\ u_x(1, 0) \\ u_x(0, 1) \\ u_x(1, 1) \\ u_y(0, 0) \\ u_y(1, 0) \\ u_y(0, 1) \\ u_y(1, 1) \\ u_{xy}(0, 0) \\ u_{xy}(1, 0) \\ u_{xy}(0, 1) \\ u_{xy}(1, 1) \end{bmatrix}$$

- Equations were hard-coded to allow for simultaneous operation on all particles

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- For simplicity, treat these as one-dimensional interpolation problems to follow two-dimensional problem
- Again, compare two methods for each: piecewise linear (faster) and piecewise cubic (more accurate)
- Linear: given $z \in [z_k, z_{k+1}]$ and 2D-interpolated values $u_k = u(x, y, z_k)$ and $u_{k+1} = u(x, y, z_{k+1})$, approximate $u(x, y, z)$ by

$$p(z) = w \cdot u_k + (1 - w) \cdot u_{k+1}$$

where

$$w = \frac{z_{k+1} - z}{z_{k+1} - z_k}$$

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- Cubic: given $z \in [z_k, z_{k+1}]$, interpolate cubic polynomial through four points $u_{k-1} = u(x, y, z_{k-1})$, $u_k = u(x, y, z_k)$, $u_{k+1} = u(x, y, z_{k+1})$, $u_{k+2} = u(x, y, z_{k+2})$
- Lagrange form:

$$\begin{aligned} p(z) = & \frac{(z - z_k)(z - z_{k+1})(z - z_{k+2})}{(z_{k-1} - z_k)(z_{k-1} - z_{k+1})(z_{k-1} - z_{k+2})} u_{k-1} \\ & + \frac{(z - z_{k-1})(z - z_{k+1})(z - z_{k+2})}{(z_k - z_{k-1})(z_k - z_{k+1})(z_k - z_{k+2})} u_k \\ & + \frac{(z - z_{k-1})(z - z_k)(z - z_{k+2})}{(z_{k+1} - z_{k-1})(z_{k+1} - z_k)(z_{k+1} - z_{k+2})} u_{k+1} \\ & + \frac{(z - z_{k-1})(z - z_k)(z - z_{k+1})}{(z_{k+2} - z_{k-1})(z_{k+2} - z_k)(z_{k+2} - z_{k+1})} u_{k+2} \end{aligned}$$

Time Integration

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- Runge-Kutta fourth order for now (simple, explicit, relatively accurate)
- Eventually will implement Milne-Hamming multistep predictor-corrector scheme in accordance with ROMS standard:

$$\hat{\mathbf{x}}_{n+1} = \mathbf{x}_{n-3} + \frac{4\Delta t}{3} (2\mathbf{u}(\mathbf{x}_n, t_n) - \mathbf{u}(\mathbf{x}_{n-1}, t_{n-1}) + \mathbf{u}(\mathbf{x}_{n-2}, t_{n-2}))$$

$$\mathbf{x}_{n+1} = \frac{9}{8}\mathbf{x}_n - \frac{1}{8}\mathbf{x}_{n-2} + \frac{3\Delta t}{8} (\mathbf{u}(\hat{\mathbf{x}}_{n+1}, t_{n+1}) + 2\mathbf{u}(\mathbf{x}_n, t_n) + \mathbf{u}(\mathbf{x}_{n-1}, t_{n-1}))$$

2D Interpolation Accuracy

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- To validate and check accuracy of 2D interpolation algorithms, applied them to analytic function $u(x, y) = e^y \sin x$ on $[0, 1]^2$
- Approximation theory in 1D:
 - n th-order interpolating polynomial p through x_0, x_1, \dots, x_n satisfies

$$u(x) - p(x) = \frac{u^{n+1}(\xi)}{(n+1)!} (x - x_0)(x - x_1) \dots (x - x_n)$$
$$\implies \|u - p\|_\infty = O(h^{n+1})$$

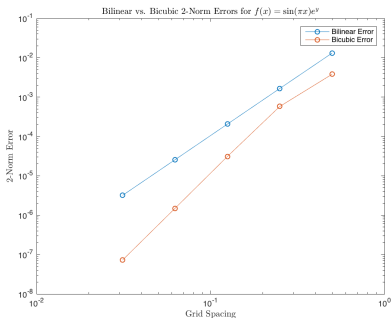
for some $\xi \in \text{int}(x_0, x_1, \dots, x_n)$, assuming $u \in C^{n+1}$

- Splines: also $O(h^{n+1})$, also assuming sufficiently high-order derivative approximations used?
- Also $\|u - p\|_\infty = O(h^{n+1})$ in 2D?

Bicubic vs. Bilinear Accuracy for $u(x, y) = e^y \sin x$

- Order of accuracy estimates ($\|u - p\|_2 = O(h^q)$):

Bilinear	Bicubic
$q \approx 3.0029$	$q \approx 4.3569$



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Bicubic vs. Bilinear Times for $u(x, y) = e^y \sin x$

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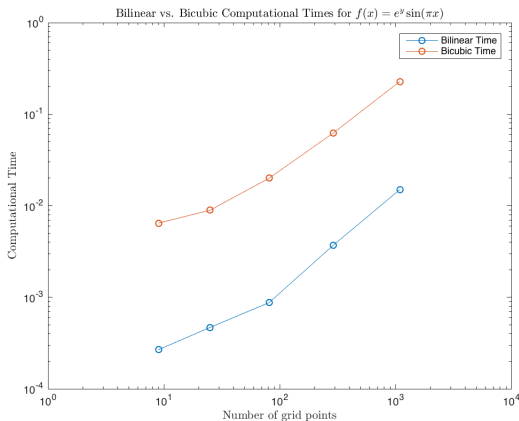
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Test Problem: Unforced, Undamped Duffing Oscillator

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- Unforced, undamped Duffing oscillator

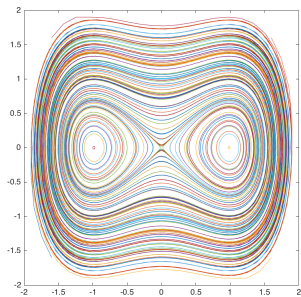
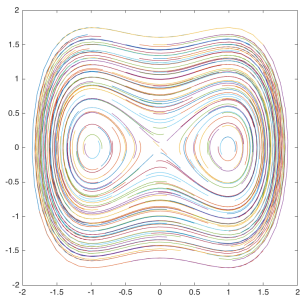
$$\frac{dx}{dt} = y$$

$$\frac{dy}{dt} = x - x^3$$

- Has exact solutions in terms of Jacobi elliptic functions

Computed Trajectories for Duffing Oscillator (RK4)

- Stable and unstable manifolds form figure-eight around fixed points
- Still need to validate against exact solution



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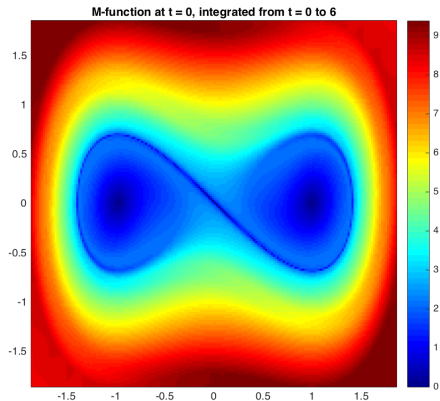
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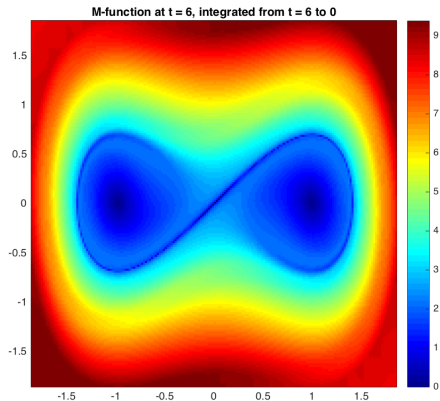
M-Function for Duffing oscillator

- Slower regions in blue, faster in red (integrated from $t = 0$ to $t = 6$)
- Blue line is *stable* manifold (particles slowing down as they approach fixed point at center)



M-Function for Duffing oscillator (integrated backwards in time)

- Slower regions in blue, faster in red (integrated from $t = 0$ to $t = 6$)
- Blue line is *unstable* manifold (particles slowing down as they approach fixed point)



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First semester:

- First half: October - Mid-November
 - Project proposal presentation and paper (done)
 - 2D and 3D interpolation
 - Need to implement cubic interpolation in vertical and time
- Second half: Mid-November - December
 - 2D trajectory implementation and validation
 - Need to validate against analytical solutions to Duffing oscillator
 - Need to validate on forced and rotating Duffing oscillator
 - M function implementation and validation
 - Mid-year report and presentation

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- Second Semester
 - First half: January - February
 - 3D trajectory implementation and validation
 - FTLE implementation
 - Second half: March - April
 - Application to ROMS dataset
 - Visualizations and further analysis
 - Final presentation and paper

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- [2] Bicubic Interpolation. (n.d.). In *Wikipedia*. Retrieved October 5, 2015 from https://en.wikipedia.org/wiki/Bicubic_interpolation.
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